

Orbits We want $r(\theta)$.

Chpt 4

Eliminate time from:

$$\text{Conservation laws } \begin{cases} \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2) + V(r) = E \\ m r^2 \dot{\theta} = J \end{cases}$$

$$\text{Combine: } \frac{1}{2} m \dot{r}^2 + \frac{J^2}{2mr^2} + V(r) = E$$

Define $u = \frac{1}{r}$. Work with this instead $u(\theta)$?

$$\frac{du}{d\theta} = -\frac{1}{r^2} \frac{dr}{d\theta}$$

$$\Rightarrow \dot{r} = \frac{dr}{d\theta} \dot{\theta} = -r^2 \dot{\theta} \frac{du}{d\theta} = -\frac{J}{m} \frac{du}{d\theta}$$

sub into radial equation

$$\frac{J^2}{2m} \left(\frac{du}{d\theta} \right)^2 + \frac{J^2}{2m} u^2 + V(u) = E$$

$V(u) = k u$ for inverse square law

As before, define length $l = \frac{J^2}{m|k|} \Rightarrow \frac{|k|}{2} l^2 \left(\frac{du}{d\theta} \right)^2 + \frac{|k|}{2} l^2 u^2 \pm k u = E$

Mult. radial eqn by $\frac{2}{|k|}$

$$l \left(\frac{du}{d\theta} \right)^2 + (u^2 \pm 2u) = \frac{2E}{|k|}$$

+ : $k > 0$

- : $k < 0$

To solve, mult. by l and add 1

$$l^2 \left(\frac{du}{d\theta} \right)^2 + l^2 u^2 \pm 2ul + 1 = \frac{2El}{|k|} + 1$$

$$\text{but } l^2 u^2 \pm 2al + 1 = (lu \pm 1)^2$$

②

Define

$$z \equiv lu \pm 1$$

$$\text{then } \frac{dz}{d\theta} = l \frac{du}{d\theta}$$

and

$$\left(\frac{dz}{d\theta}\right)^2 + z^2 = \frac{2El}{|k|} + 1 \equiv e^2$$

dimensionless constant

N.B. only has solⁿ when $\frac{2El}{|k|} + 1 > 0$

$$\text{i.e. } E > -\frac{|k|}{2l}$$

General solution of equation: $z = e \cos(\theta - \theta_0)$

\uparrow const. of integration

but $z = l \frac{1}{r} \pm 1$, so

$$r[e \cos(\theta - \theta_0) - 1] = l$$

$$k > 0$$

repulsive

$$r[e \cos(\theta - \theta_0) + 1] = l$$

$$k < 0$$

attractive

\Rightarrow Conic sections with one focus as origin

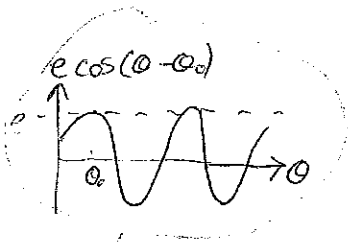
e = eccentricity

l = semi-latus rectum

θ_0 = orientation of l wrt. coordinate axes

N.B. repulsive case: e must be ≥ 1 . But $e^2 = \frac{2El}{|k|} + 1$ so $E > 0$
(else [] always -ve)

attractive case: when $e \geq 1 \Rightarrow E > 0$, then $r \rightarrow \infty$ when $[] = 0$
 i.e. particle escapes if $E > 0$



For both cases:

(3)

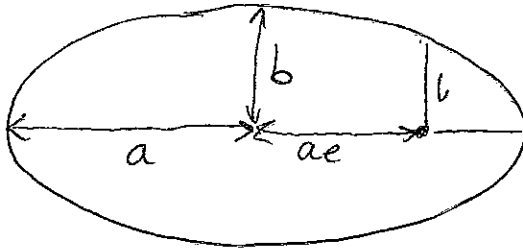
r takes its min value when $\theta = \theta_0$

$\Rightarrow \theta_0$ specifies direction of pt of closest approach

In attractive case l has ^{simple} geometric meaning: $r=l$ when $\theta = \theta_0 \pm \frac{\pi}{2}$

$\Rightarrow l$ is radial distance at right angles to closest approach

Elliptic orbits ($E < 0, 0 \leq e < 1$)



Re-write orbit equation in cartesian form, choose $\theta_0 = 0$

$$r(e \cos \theta \pm 1) = l \quad \begin{array}{l} \nearrow \text{attractive} \\ \searrow \text{repulsive} \end{array} \quad \begin{array}{l} e \geq 0 \\ l > 0 \end{array}$$

Write as
$$\underbrace{r \cos \theta}_{e x} \pm r = l$$

or
$$\pm r = l - e x$$

Square:
$$r^2 = x^2 + y^2 = (l - e x)^2$$

or
$$x^2 - e^2 x^2 + 2l e x + y^2 = l^2$$

$$\Rightarrow (1 - e^2) x^2 + 2l e x + \frac{e^2 l^2}{1 - e^2} - \frac{e^2 l^2}{1 - e^2} + y^2 = l^2$$

$$\Rightarrow \left(\sqrt{1 - e^2} x + \frac{e l}{\sqrt{1 - e^2}} \right)^2 + y^2 = l^2 + \frac{e^2 l^2}{1 - e^2} = l^2 \left(1 + \frac{e^2}{1 - e^2} \right)$$

$$\left(\sqrt{1-e^2} x + \frac{el}{\sqrt{1-e^2}} \right)^2 + y^2 = \frac{l^2}{1-e^2} \quad (4)$$

or $(1-e^2) \left(x + \frac{el}{1-e^2} \right)^2 + y^2 = \frac{l^2}{1-e^2}$

$$\Rightarrow \frac{(1-e^2)^2}{l^2} \left(x + \frac{el}{1-e^2} \right)^2 + y^2 \frac{(1-e^2)}{l^2} = 1$$

Define $a \equiv \frac{l}{1-e^2}$, $b \equiv \frac{l}{\sqrt{1-e^2}}$ semi-axes

then $\frac{(x+ea)^2}{a^2} + \frac{y^2}{b^2} = 1$

same as $\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$ but with origin at $(-ae, 0)$.

N.B. $e^2 = \frac{a^2 - b^2}{a^2}$ $l = \frac{b^2}{a}$

Physically $a = \frac{l}{1-e^2} = \frac{\cancel{m} l}{1 - \frac{(1+2E/l)}{\cancel{m} l}} = \frac{\cancel{m} l}{-2mE} = \frac{l}{2|E|}$

semi-axis a is determined by energy E .

$b^2 = al = \frac{l}{2|E|} \frac{J^2}{m l} = \frac{J^2}{2m|E|}$ So l is fixed by J .

Time taken to traverse any part of orbit is found (5)

from $\frac{dA}{dt} = \frac{J}{2m}$

$$\Rightarrow dt = \frac{dA}{\frac{J}{2m}}$$

↑
area swept off by radius vector

Period of orbit? $A = \pi ab$

$$\text{so } T = \pi ab \frac{2m}{J}$$

$$\text{or } \left(\frac{T}{2\pi}\right)^2 = \frac{m^2 a^2 b^2}{J^2} = \frac{m^2 a^3 l}{\cancel{J^2}}$$

$$\text{but } l = \frac{J^2}{m|k|}$$

$$\text{so } \left(\frac{T}{2\pi}\right)^2 = \frac{m^2 a^3 \cancel{J^2}}{m|k| \cancel{J^2}} = \frac{m}{|k|} a^3$$

So for gravity $k = GM$ and $\left(\frac{T}{2\pi}\right)^2 = \frac{a^3}{GM}$

Kepler's 3rd law.

Hyperbolic orbits ($E > 0, e > 1$)

Again $x^2 + y^2 = (l - e \cos \theta)^2$

Complete square as before:

$$(1 - e^2)x^2 + 2elx + \frac{e^2 l^2}{1 - e^2} - \frac{e^2 l^2}{1 - e^2} + y^2 = l^2$$

multi. by -1:

$$(e^2 - 1)x^2 - 2elx + \frac{e^2 l^2}{e^2 - 1} - \frac{e^2 l^2}{e^2 - 1} - y^2 = -l^2$$

$$\left(\sqrt{e^2 - 1} x - \frac{el}{\sqrt{e^2 - 1}} \right)^2 - y^2 = \frac{e^2 l^2}{e^2 - 1} - l^2 = l^2 \left(\frac{e^2}{e^2 - 1} - 1 \right) = \frac{l^2}{e^2 - 1}$$

$$\Rightarrow (e^2 - 1) \left(x - \frac{el}{e^2 - 1} \right)^2 - y^2 = \frac{l^2}{e^2 - 1}$$

$$\Rightarrow \frac{(e^2 - 1)^2}{l^2} \left(x - \frac{el}{e^2 - 1} \right)^2 - y^2 \frac{(e^2 - 1)}{l^2} = 1$$

Define $a = \frac{l}{e^2 - 1}$, $b = \frac{l}{\sqrt{e^2 - 1}}$

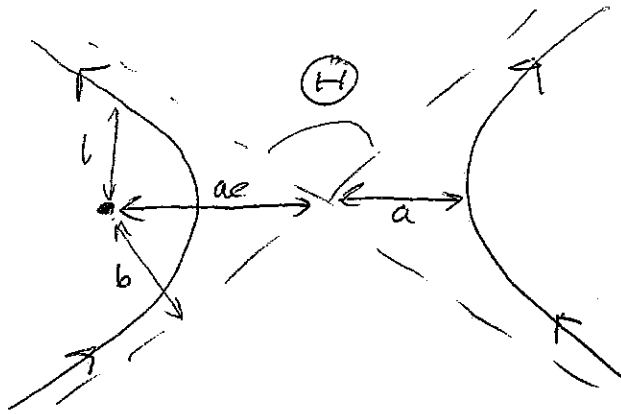
$$\frac{(x - ea)^2}{a^2} - \frac{y^2}{b^2} = 1$$

\Rightarrow hyperbola with centre at $(ae, 0)$

left hand branch = orbit under attractive $1/r^2$ force (intersects x axis at $(ae - a, 0)$)

right " " = " " repulsive (" " " " $(ae + a, 0)$)

(7)

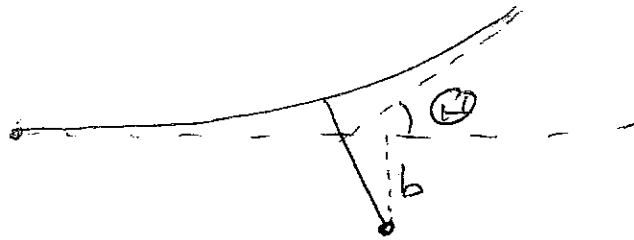


note that $a = \frac{l}{e^2 - 1} = \frac{|l|}{2E}$

$$b^2 = a l = \frac{J^2}{2mE}$$

a is determined by energy
 l " " " J

A to semiminor axis b is identical to impact parameter



Asymptotes :

$$\Theta = \pm \arccos\left(\frac{1}{e}\right)$$

repulsive

i.e. $r \rightarrow \infty$

i.e. $(e \cos \Theta \pm 1) = 0$
 from $r(e \cos \Theta \pm 1) = b$
attractive
rep

$$\Theta = \pm [\pi - \arccos\left(\frac{1}{e}\right)]$$

attractive

Deflection angle in both cases:

scattering angle

$$\Theta = \pi - 2 \arccos\left(\frac{1}{e}\right)$$

What is relation between impact parameter b, and scattering angle Θ , and limiting velocity v at infinity?

(8)

from $\Theta = \pi - 2 \arccos\left(\frac{1}{e}\right)$

$$\Rightarrow e = \sec\left[\frac{1}{2}(\pi - \Theta)\right] \quad \sec = \frac{1}{\cos}$$

sub in $b^2 = al = \frac{l^2}{e^2 - 1} = a^2(e^2 - 1)$ $\left[a^2 = \frac{l}{e^2 - 1}\right]$

$$\begin{aligned} &= a^2 \left[\sec^2\left[\frac{1}{2}(\pi - \Theta)\right] - 1 \right] = a^2 \left[\frac{1}{\cos^2} - \frac{\cos^2}{\cos^2} \right] \\ &= a^2 \cot^2\left(\frac{\Theta}{2}\right) \end{aligned} \quad \begin{aligned} &= a^2 \frac{1 - \cos^2}{\cos^2} = a^2 \frac{\sin^2\left(\frac{\pi - \Theta}{2}\right)}{\cos^2\left(\frac{\pi - \Theta}{2}\right)} \\ &= a^2 \frac{\cot^2\left(\frac{\Theta}{2}\right)}{\sin^2\left(\frac{\Theta}{2}\right)} \end{aligned}$$

but also, $a = \frac{|k|}{2E} = \frac{|k|}{2 \frac{1}{2} m v^2} = \frac{|k|}{m v^2}$

$$\Rightarrow b = \frac{|k|}{m v^2} \cos\left(\frac{\Theta}{2}\right)$$